

# Japan MEXT Leading Project for Laser-Produced Plasma EUV Light Source Development

**Katsunobu Nishihara**  
**Institute of Laser Engineering, Osaka University**

## Outline of talk

- requirement of light source and design window
- laser intensity scaling of the conversion efficiency  
experiments and modeling (Sn, SnO<sub>2</sub>)  
opacity and satellite line effect (Xe)
- further improvement of conversion efficiency (Sn, SnO<sub>2</sub>)

## In collaboration with

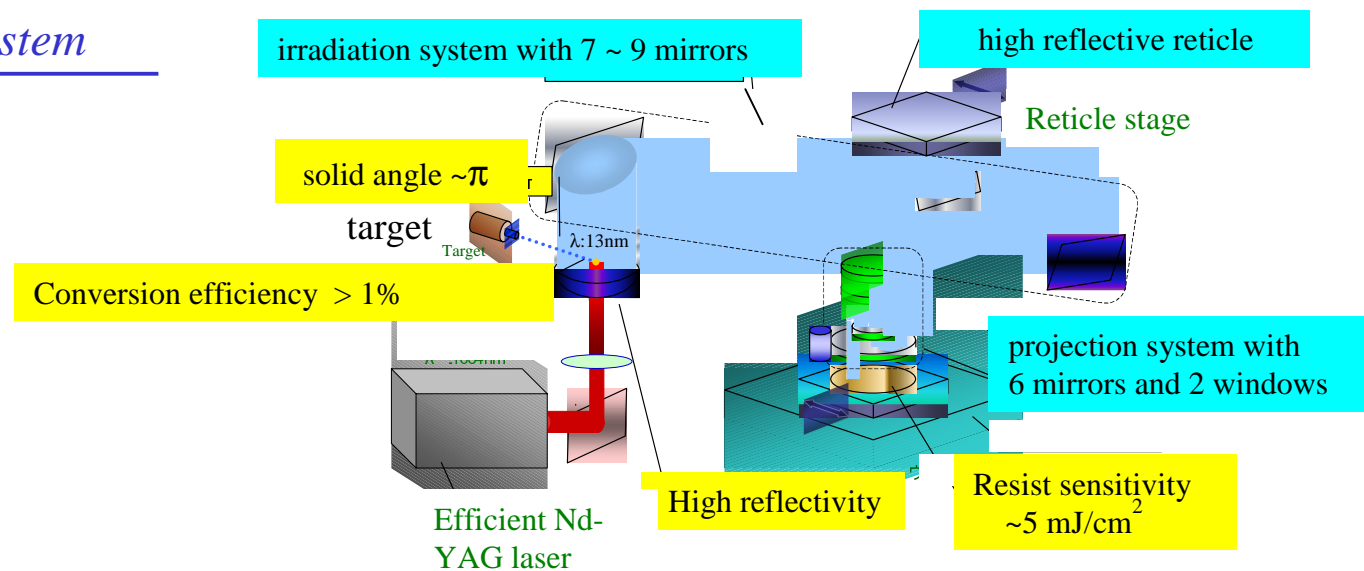
(theory) : T. Nishikawa, A. Sasaki, T. Kawamura, A. Sunahara, H. Furukawa,  
M. Murakami, R. More, K. Fujima, F. Koike, T. Kagawa, V. Zhakhovskii,  
T. Kato and H. Tanuma,

(experiments) : H. Nishimura, M. Nakai, K. Shigemori, Y. Shimada, S. Uchida,  
S. Fujioka, K. Hashimoto, T. Hibino, R. Matsui, T. Okuno, F. L. Sohnatzadeh,  
Y. Tao, M. Yamaura, K. Nagai, T. Norimatsu, M. Nakatsuka, H. Fujita,  
Y. Fujimoto, K. Tsubakimoto, H. Yoshida, N. Miyanaga and Y. Izawa

**A new-five-year-long project for LPP EUV light source development has been started as the Leading Project promoted by MEXT.**

**The project aims at understanding physics of EUV light source for lithography, and providing database and guideline for practical use in collaboration with METI project (EUVA).**

*EUVL system*



**The project includes high performance laser development ( 5kW, 5kHz), plasma experiments, theory and simulation, and target fabrication.**

## Requirements of the light source and possible design windows

wavelength: 13.5 nm @2%BW / EUV power: > 115 W (after intermediate focus)  
 repetition: 10 kHz / etendue: 1 mm<sup>2</sup>str / conversion efficiency: > 1 %

**EUV power at source / repetition :** 300 W / 10 kHz  
**EUV energy per shot :** 30 mJ / shot  
**source plasma size (  $S \pi = \pi^2 r^2 = 1 \text{ mm}^2 \text{str}$  ) :**  $\Phi \approx 600 \mu\text{m}$   
**EUV intensity (  $I_{\text{EUV}} = E_{\text{EUV}} / (S \tau_{\text{EUV}})$  ) :**  $10^9 \sim 10^{10} \text{ W/cm}^2$   
 ( pulse width :  $\tau_{\text{EUV}} = 1 \sim 10 \text{ ns}$  )  
**laser intensity (  $I_{\text{L}} = I_{\text{EUV}} / \eta_{\text{EUV}}$  ) :**  $10^{11} \sim 10^{12} \text{ W/cm}^2$   
 ( conversion efficiency :  $\eta_{\text{EUV}} = 0.01$  )  
**ion density ( Xe :  $n_{5p-4d}$ , Sn :  $n_{4f-4d}$  ) :**  $10^{18} \sim 10^{20} \text{ cm}^{-3}$   
 ( photon number per unit time and area :  
 $E_{\text{EUV}} / (h\nu \tau_{\text{EUV}} S) = 4 l_{\text{OD}=1} n_{4f-4d \text{ or } 5p-4d} A_{\text{Eien}}$   
**opacity :  $l_{\text{OD}=1} = 10 \mu\text{m}$ , oscillator strength :  $g_n / g_n f_{n'n} = 1$ ,**  
**abundance of  $n_{4f-4d \text{ or } 5p-4d} = 0.1$  )**

KN030806-0



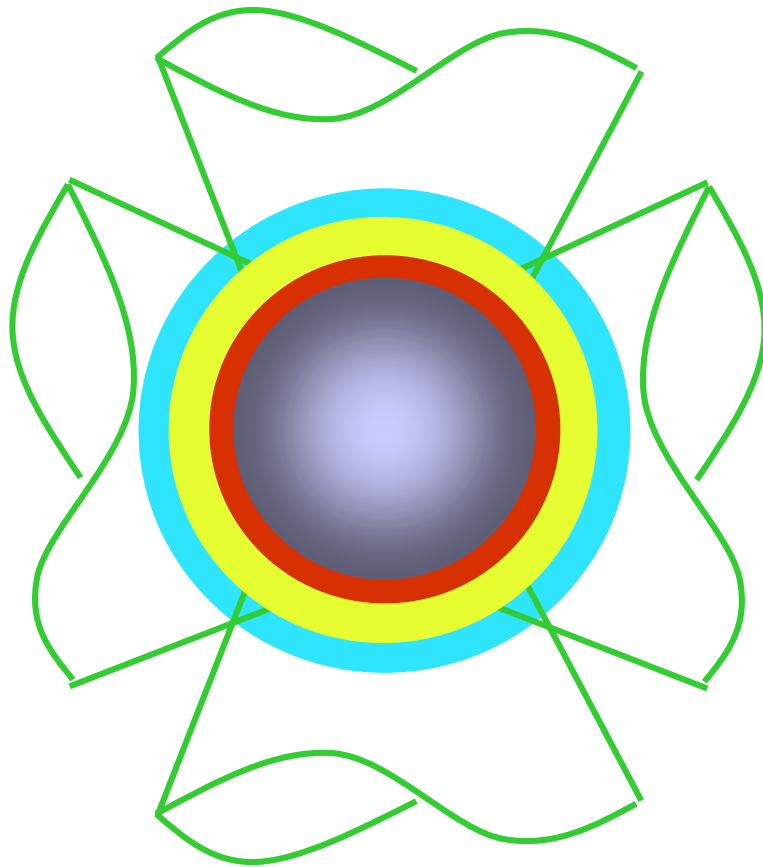
## Executive summary

- **The conversion efficiency of a few % to 13.5 nm EUV light of 2 % bandwidth has been attained at  $1 - 2 \times 10^{11}$  W/cm<sup>2</sup>. (Sn)**
- **Considering power balance, we present a scaling law of the conversion efficiency, which agrees fairly well the experiments.**
- **We show that 13.5 nm EUV light is emitted in both the high-temperature corona and the high-density region, depending on the laser intensity.**
- **We also show that the Planckian intensity can be achieved for the lines due to the effects opacity and satellite lines in the high-density region.**
- **$2\omega_0$  laser irradiation on SnO<sub>2</sub> target leads to further improvement of the conversion efficiency at intensity of  $.5 - 1 \times 10^{11}$  W/cm<sup>2</sup>, although further study is required.**

KN030930-1



An ideally “uniform” EUV radiator was produced by GEKKO-XII laser, to obtain laser intensity dependence of the conversion efficiency without lateral energy loss and geometrical effects.



spherically uniform plasmas

**Laser :**

**GEKKOXII, 12 beams**  
**wavelength:  $\omega$  (1.056  $\mu\text{m}$ ),  $2\omega$  (0.527  $\mu\text{m}$ )**  
**intensity:  $5 \times 10^{10} \sim 1 \times 10^{12} \text{ W/cm}^2$**   
**pulse width: 1.2 ns (FWHM, Gaussian)**

**Target :**

**Sn or SnO<sub>2</sub> coated glass ball**  
**300~2000  $\mu\text{m}^\phi$  ( mostly 700  $\mu\text{m}^\phi$ )**

**Diagnostics (XST: time resolved):**

**transmission grating (TDI) + CCD**  
**grazing incident spectrometers (GIS)**

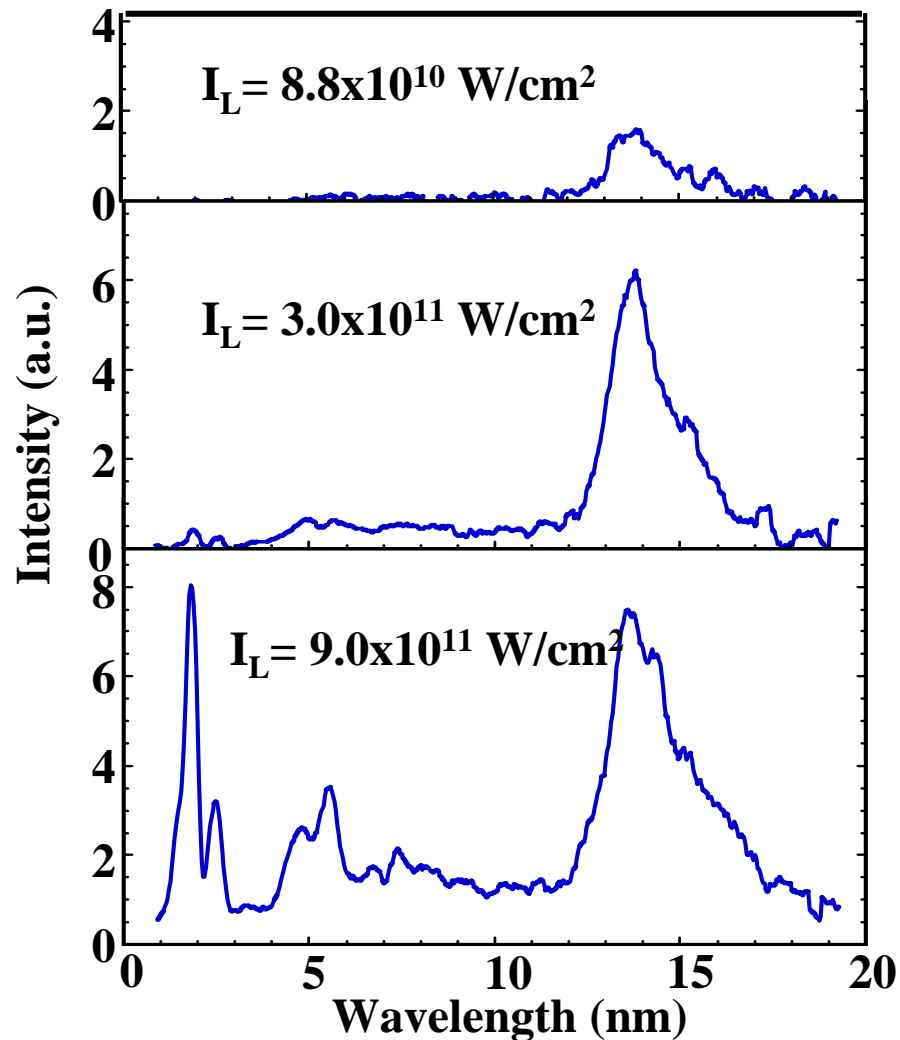
**x-ray backlight: density profile**



See Nakai(101), Nishimura(102), Yamamura(92)  
posters for detail experimental results



**Laser intensity dependence of spectra shows that short wavelength (1-3nm) emission appears at high intensity ( $10^{12}\text{W}/\text{cm}^2$ ), while long wavelength ( $>10\text{nm}$ ) emission dominates at low intensity ( $10^{11}\text{W}/\text{cm}^2$ ).**



target Sn, laser  $\omega$   
transmission grating + CCD

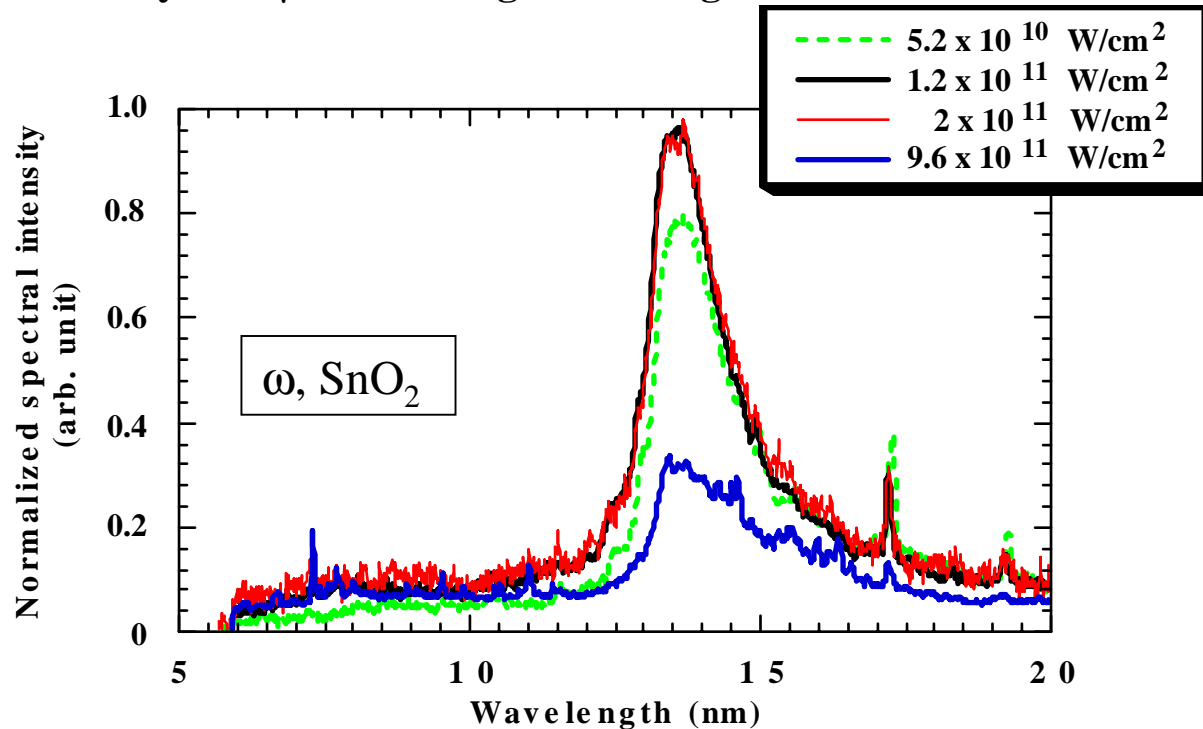
- Emission with wavelength of 13-15 nm is mainly due to 4d-4f transition.
- M-shell emission causes short wavelength emission  $< 3\text{nm}$  (1 - 0.5 keV), while N-shell emission does longer wavelength emission  $> 4\text{nm}$  ( $< 0.5 \text{ keV}$ ).

KN-YS030806-1

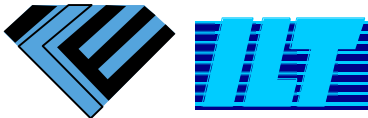


Spectral intensity normalized by laser intensity shows that conversion efficiency has its maximum around  $1-2 \times 10^{11} \text{ W/cm}^2$ .

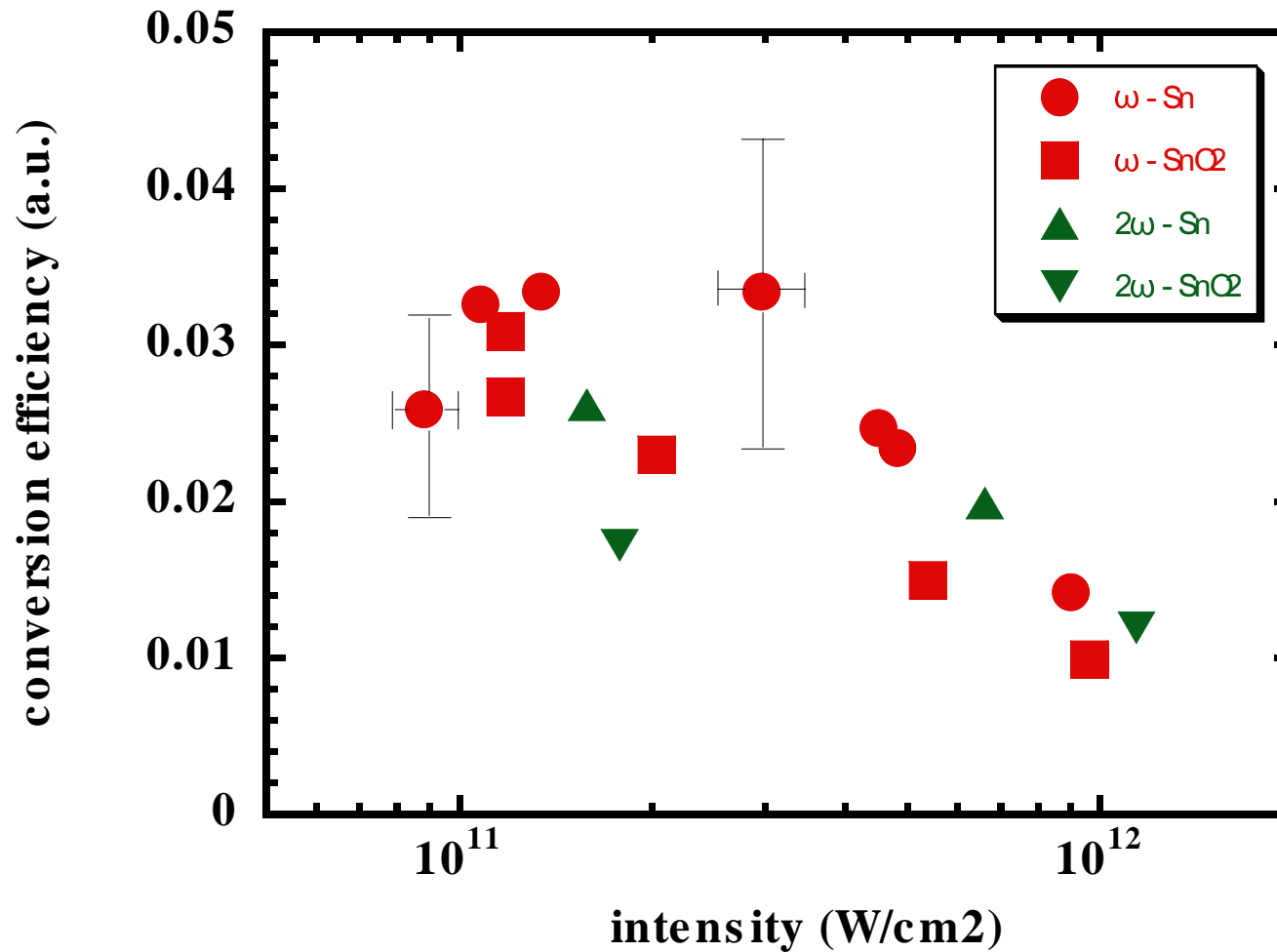
Intensity dependence of EUV spectrum from spherical tin-oxide targets illuminated by  $1.05\text{-}\mu\text{m}$  wavelength laser light.



Detailed spectral features are observed by using grazing incident spectrometers (GIS) coupled with a back illumination CCD camera and an x-ray streak camera.



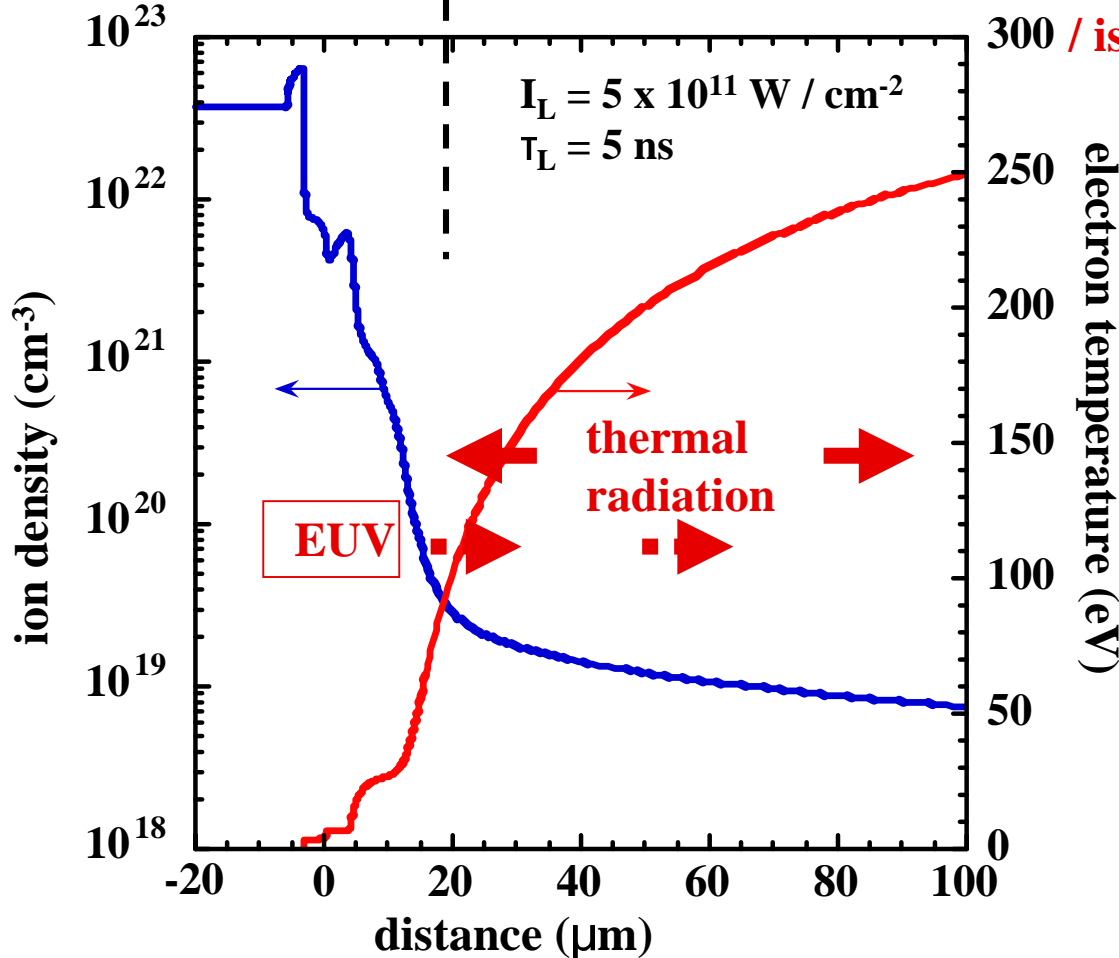
Transmission grating data also shows high conversion efficiency near  $1\text{-}2 \times 10^{11}$  W/cm<sup>2</sup> and relatively weak laser wavelength dependence.



# Features of Laser Produced High-Z Plasma

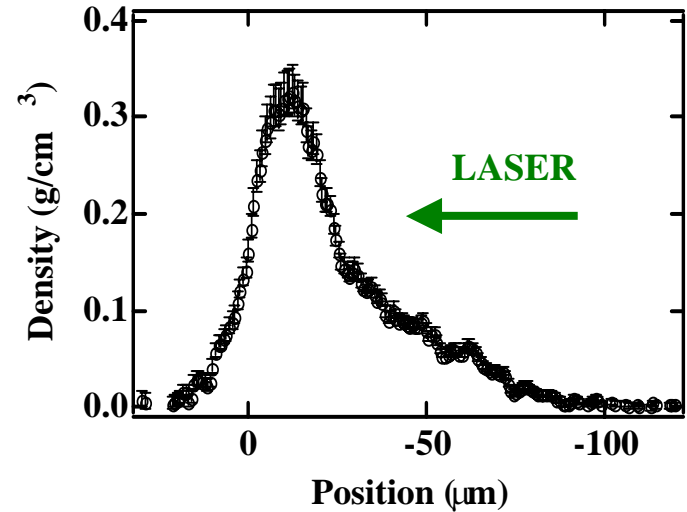
High density : high density / low temperature / LTE / b-f absorption region

Corona region : high temperature / low density / CRE / M-, N- band emission / isothermal expansion



High density region was observed by side-on x-ray back light method .

$I_L = 1.2 \times 10^{12} \text{ W / cm}^{-2}$



KN030806-1

See Sunahara's(85) posters for details



## Modeling of EUV-LPP (power balance)

Consider following energy losses

- kinetic energy loss due to expansion,  $E_k$
- ionization loss,  $E_i$
- radiation loss,  $E_R$ ,

and assume a sum of the losses equals to laser flux.

$$I_L = E_k + E_i + E_R$$

Assume **isothermal expansion in corona**

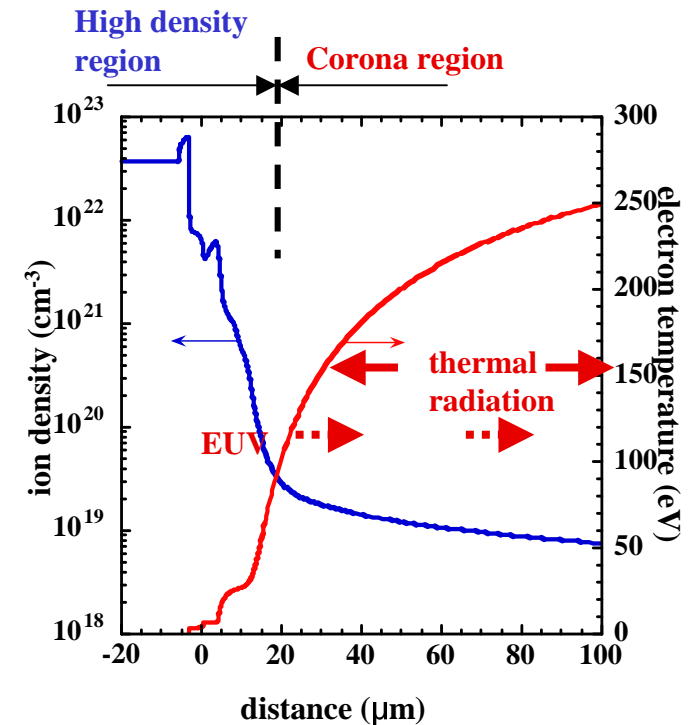
$$n = n_0 \exp(-x/c_s t), \quad v = c_s + x/t$$

$$E_k = 3 Z^*(n_i, T_e) n_i T_e c_s,$$

$$E_i = E_i(n_i, T_e, Z^*(n_i, T_e)) n_i c_s,$$

$$E_R = P_R(n_i, T_e) c_s \tau_L / \alpha,$$

where  $Z^*$  : ionization state,  $n_i$  : ion density,  
 $T_e$  : electron temperature,  $c_s$  : sound speed  
 $\tau_L$  : laser pulse duration,  $\alpha (\approx 1.5) : P_R \propto n_i^\alpha$ ,

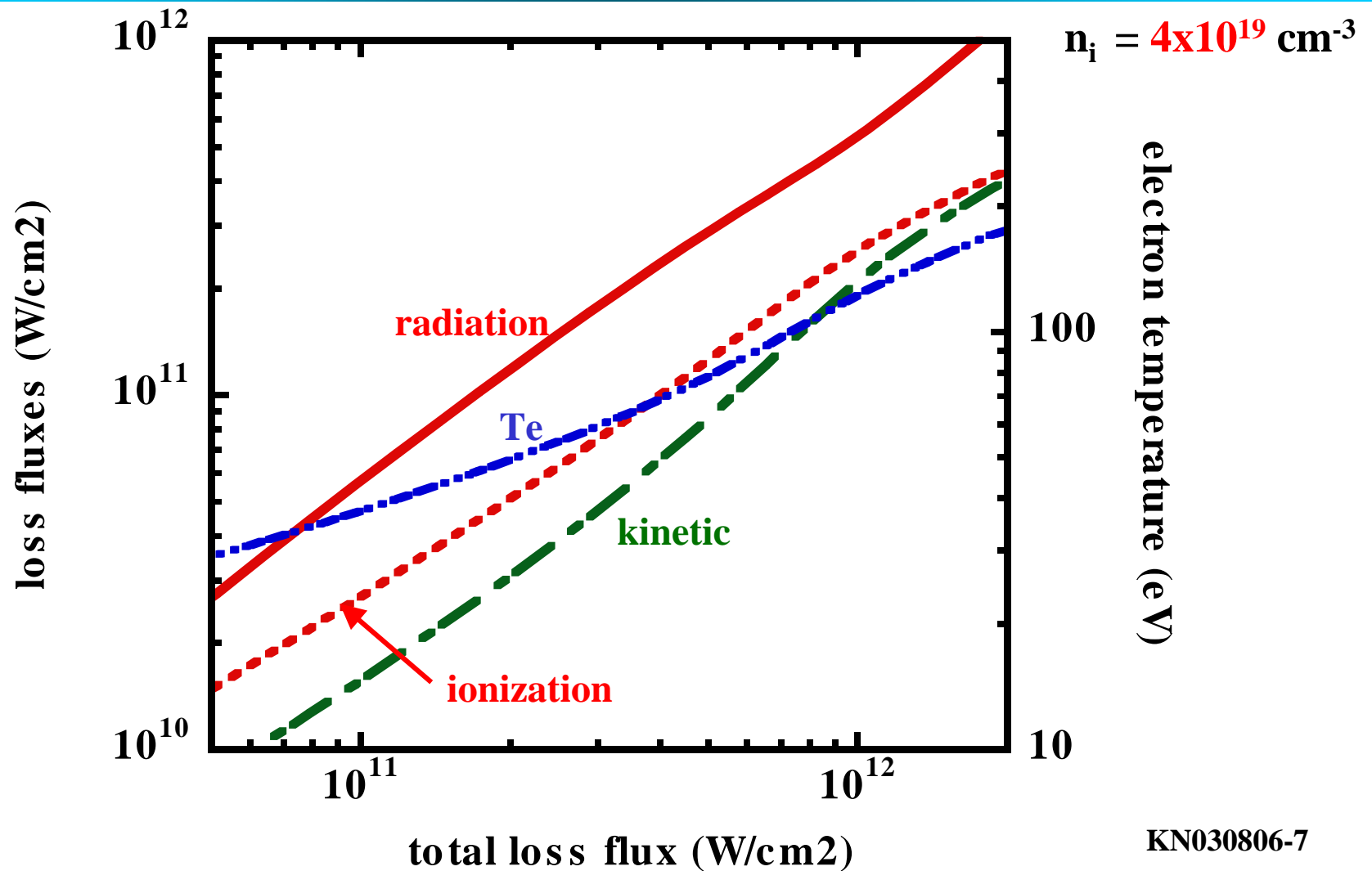


KN030806-2

See Nishikawa(81) and Sunahara's(85) posters for details



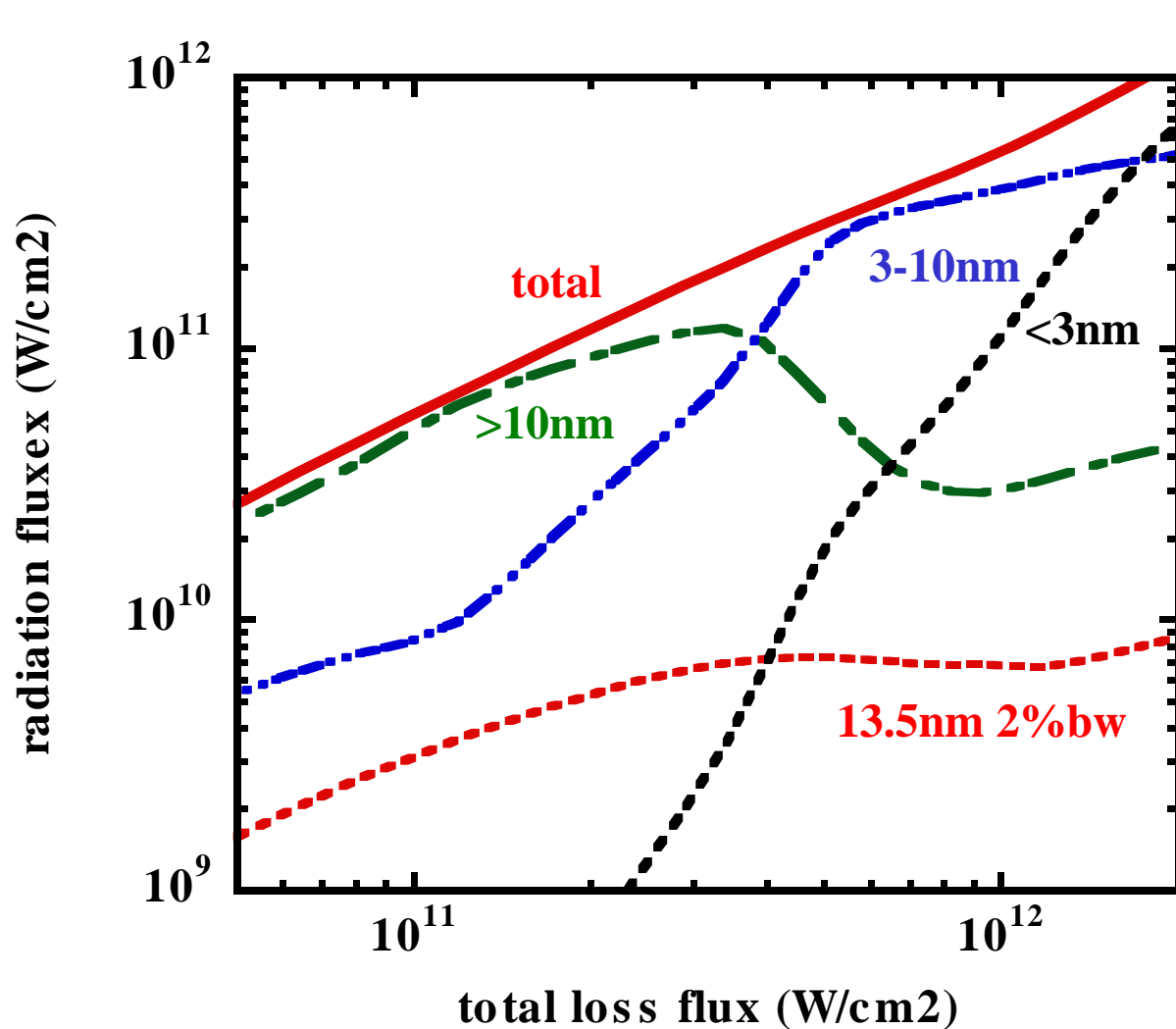
Dependence of the losses on laser intensity indicates that radiation loss is dominant (50%), ionization loss is large at low intensity and kinetic loss becomes large at high intensity, but still small.



KN030806-7



Model indicates that radiation spectra may drastically change with laser intensity, which agrees well with the experiments.



$$n_i = 4 \times 10^{19} \text{ cm}^{-3}$$

13.5 nm with 2% bandwidth assumptions

- 10% of 12.4-14.5 nm emission from the corona calculated from AAM with  $(n,l)$  splitting
- Planck radiation from the high-density region with radiation temperature estimated from the total radiation from the corona.

KN030806-9



**Both opacity and satellite lines play important roles in LPP.  
( Xe plasma with 4d-4f, 4p-4d and 4d-5p satellite lines )**

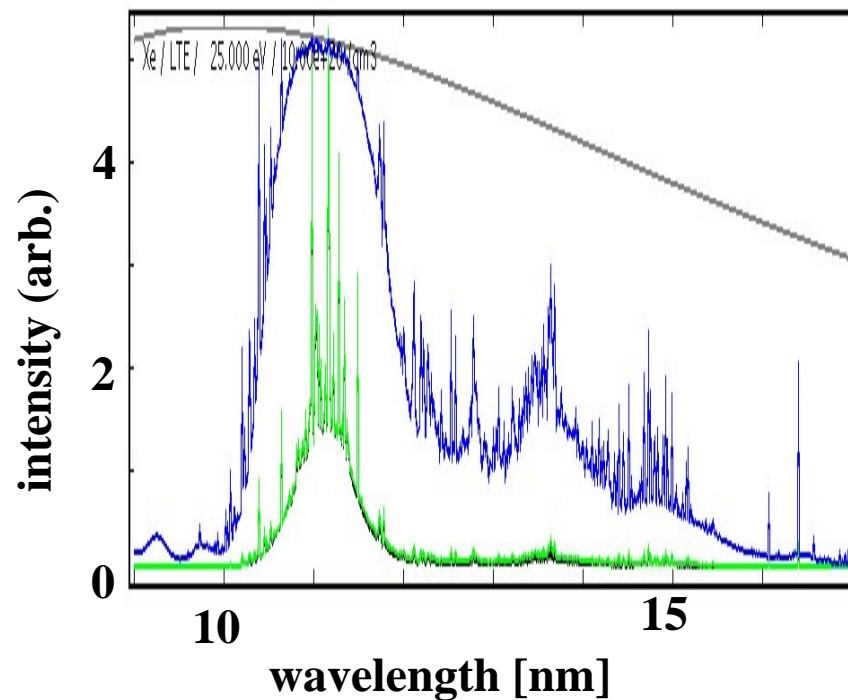
**Plasma parameters**

$$T_e = 25\text{eV}$$

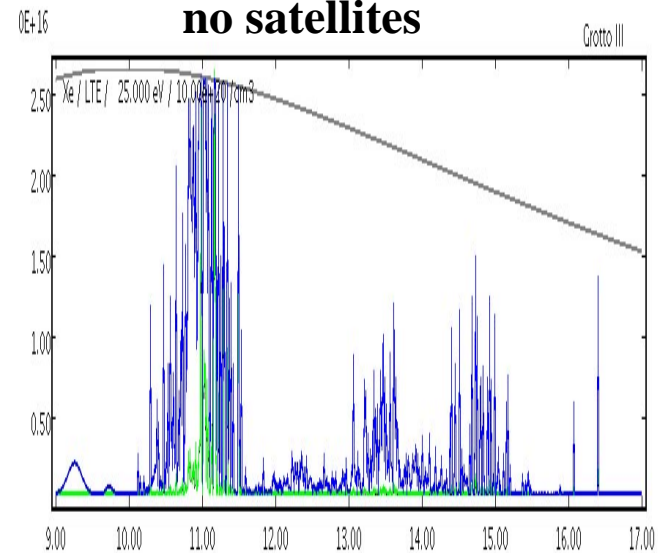
$$n_e = 10^{21}/\text{cm}^3$$

$$r = 10\mu\text{m}$$

**with satellites**



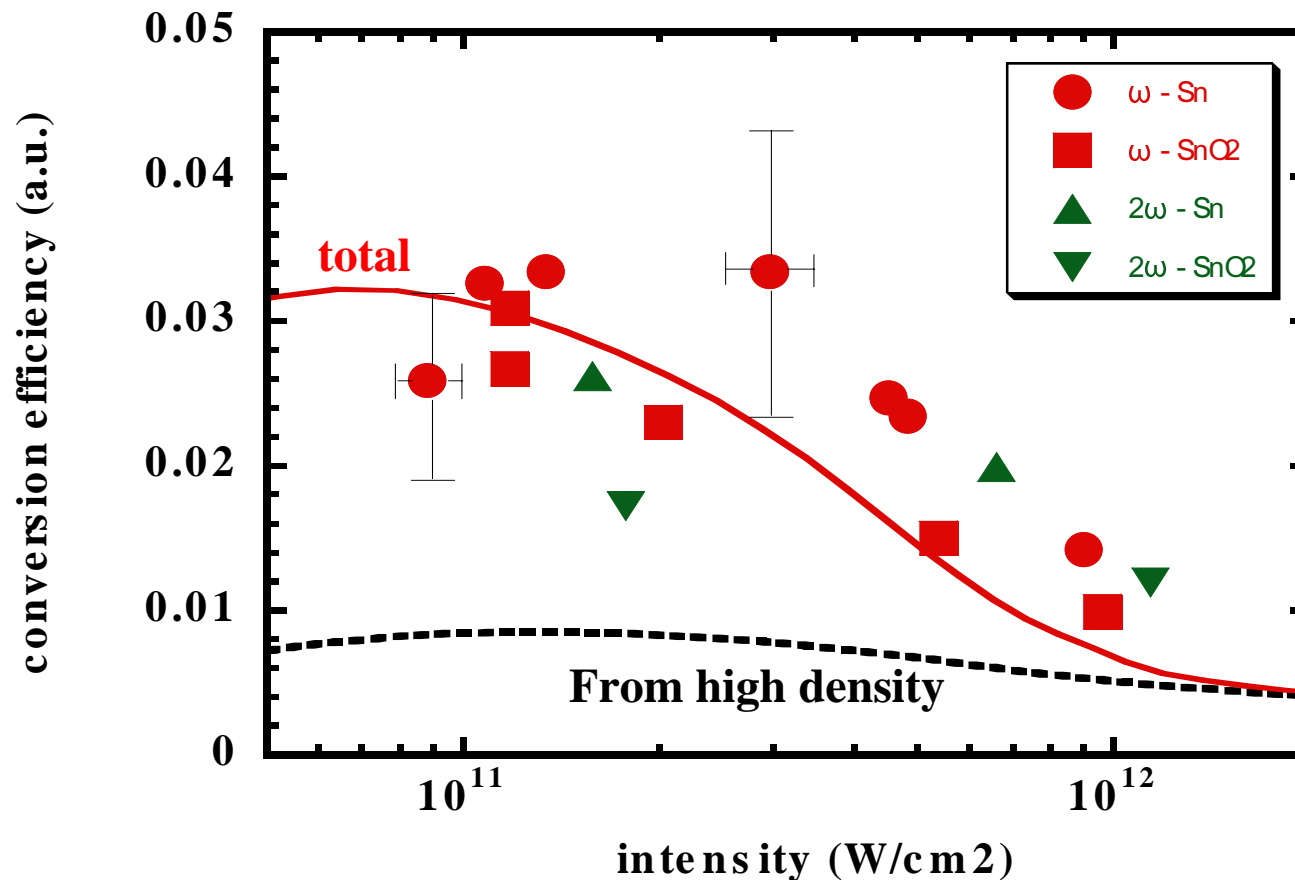
**no satellites**



**See Sasaki's(79) posters for details**

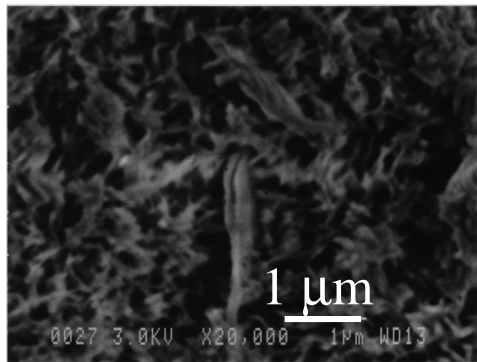


Theoretical conversion efficiency agrees fairly well with the experiments. At low intensity, emission from the corona is dominate, while that from the high density region becomes large at high intensity.

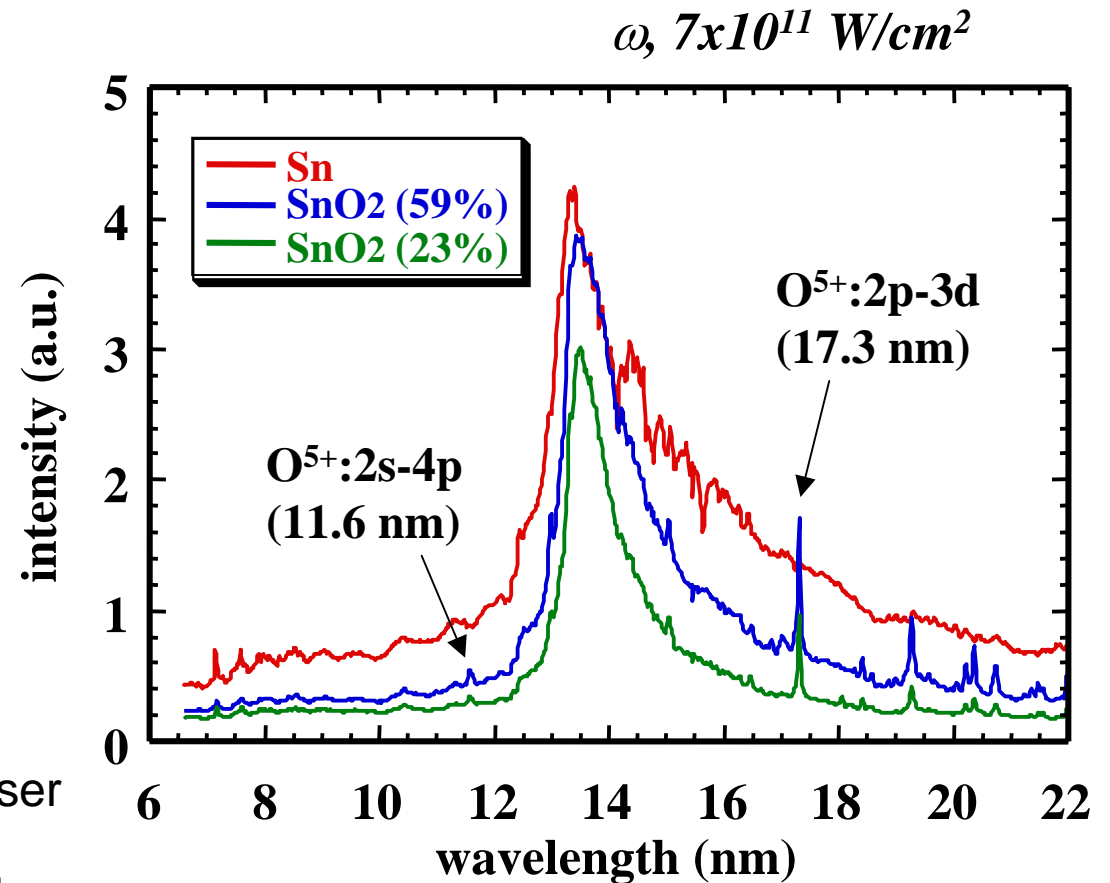
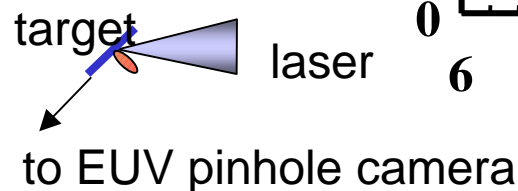
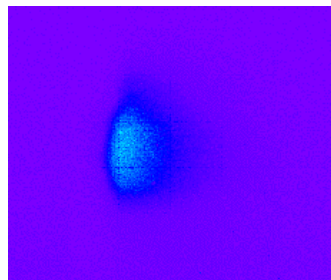


With decrease in density, spectrum becomes narrower but emission at 13.5 nm does not largely change. (planar target)

low density SnO<sub>2</sub>  
(23% of solid density)



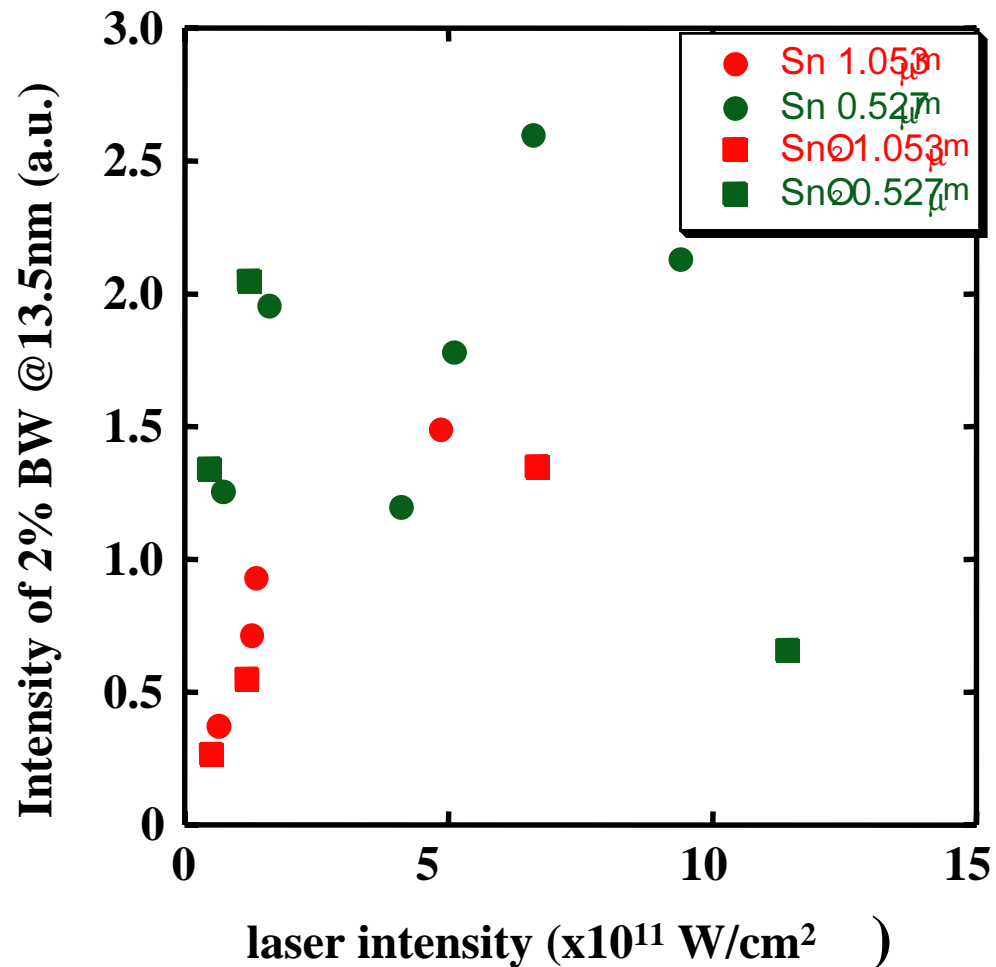
pinhole image  
(@13.5 nm 2%BW)



See Nishimura(102), Nagai(93)  
posters for detail experimental results



At low intensity of  $.5 - 1 \times 10^{11} \text{ W/cm}^2$ ,  $2\omega$  laser irradiation results in higher conversion efficiency, and  $\text{SnO}_2$  is higher Sn (?).



See Yamaura(92) posters for detail experimental results



## Conclusions

---

- **The conversion efficiency of a few % to 13.5 nm EUV light of 2 % bandwidth has been attained at 1 -  $2 \times 10^{11}$  W/cm<sup>2</sup>.**
- **Considering power balance, we present a scaling law of the conversion efficiency, which agrees fairly well the experiments.**
- **We show that 13.5 nm EUV light is emitted in both the high-temperature corona and the high-density region, depending on the laser intensity.**
- **We also show that the Planckian intensity can be achieved for the lines due to the effects opacity and satellite lines in the high-density region.**
- **$2\omega_0$  laser irradiation on SnO<sub>2</sub> target leads to further improvement of the conversion efficiency at intensity of  $.5-1 \times 10^{11}$  W/cm<sup>2</sup>, although further study is required.**

KN030930-1

---

**Those results indicate that the source parameters for practical use can be achieved.**

